

Notes, books, and calculators are not authorized. Show all your work in the blank space you are given on the exam sheet. **Answers with no justification will not be graded.**

Here are some formulae that you may want to use:

$$\mathcal{F}(f)(\omega) \stackrel{\text{def}}{=} \frac{1}{2\pi} \int_{-\infty}^{+\infty} f(x)e^{i\omega x} dx, \quad \mathcal{F}^{-1}(f)(x) = \int_{-\infty}^{+\infty} f(\omega)e^{-i\omega x} d\omega, \quad (1)$$

$$\mathcal{F}(f(x - \beta))(\omega) = e^{i\beta\omega} \mathcal{F}(f)(\omega), \quad (2)$$

$$\mathcal{F}(f * g) = 2\pi \mathcal{F}(f)\mathcal{F}(g), \quad (3)$$

$$\mathcal{F}(e^{-\alpha x^2}) = \frac{1}{\sqrt{4\pi\alpha}} e^{-\frac{\omega^2}{4\alpha}} \quad (4)$$

Question 1: Consider the equation $\partial_x \left(\frac{1}{1+3x^2} \partial_x u(x) \right) = f(x)$, $x \in (0, 1)$, $\partial_x u(0) = a$, $u(1) = b$. Let $G(x, x_0)$ be the associated Green's function.

(i) Give the equation and boundary conditions satisfied by G .

The operator is clearly self-adjoint. Then for all $x \neq x_0$ we have

$$\partial_x \left(\frac{1}{1+3x^2} \partial_x G(x, x_0) \right) = \delta_{x-x_0}, \quad \partial_x G(0, x_0) = 0, \quad G(1, x_0) = 0.$$

(ii) Give the integral representation of $u(x_0)$ for all $x_0 \in (0, 1)$ in terms of G , f , and the boundary data. (Do not compute G in this question).

Multiply the equation defining G by u and integrate over $(0, 1)$,

$$\langle \delta_{x-x_0}, u \rangle = u(x_0) = \int_0^1 \partial_x \left(\frac{1}{1+3x^2} \partial_x G(x, x_0) \right) u(x) dx.$$

We integrate by parts and we obtain

$$\begin{aligned} u(x_0) &= - \int_0^1 \frac{1}{1+3x^2} \partial_x G(x, x_0) \partial_x u(x) dx + \left[\frac{1}{1+3x^2} \partial_x G(x, x_0) u(x) \right]_0^1 \\ &= \int_0^1 G(x, x_0) \partial_x \left(\frac{1}{1+3x^2} \partial_x u(x) \right) dx + \frac{1}{4} \partial_x G(1, x_0) u(1) - \left[G(x, x_0) \frac{1}{1+3x^2} \partial_x u(x) \right]_0^1 \\ &= \int_0^1 G(x, x_0) \partial_x \left(\frac{1}{1+3x^2} \partial_x u(x) \right) dx + \frac{1}{4} \partial_x G(1, x_0) u(1) + G(0, x_0) \partial_x u(0). \end{aligned}$$

Now, using the boundary conditions and the fact that $\partial_x((1+3x^2)^{-1} \partial_x u(x)) = f(x)$, we finally have

$$u(x_0) = \int_0^1 G(x, x_0) f(x) dx + \frac{1}{4} \partial_x G(1, x_0) b + G(0, x_0) a.$$

(iii) Compute $G(x, x_0)$ for all $x, x_0 \in (0, 1)$.

For all $x \neq x_0$ we have

$$\partial_x \left(\frac{1}{1+3x^2} \partial_x G(x, x_0) \right) = \delta_{x-x_0}, \quad \partial_x G(0, x_0) = 0, \quad G(1, x_0) = 0.$$

The generic solution is

$$G(x, x_0) = \begin{cases} a(x+x^3) + b & \text{if } 0 \leq x < x_0 \\ c(x+x^3) + d & \text{if } x_0 < x \leq 1. \end{cases}$$

The boundary conditions give

$$\partial_x G(0, x_0) = 0 = a, \quad G(1, x_0) = 0 = 2c + d.$$

As a result

$$G(x, x_0) = \begin{cases} b & \text{if } 0 \leq x < x_0 \\ c(x+x^3) - 2c & \text{if } x_0 < x \leq 1. \end{cases}$$

G must be continuous at x_0 ,

$$b = c(x_0 + x_0^3) - 2c$$

and must satisfy the gap condition

$$\int_{x_0-\epsilon}^{x_0+\epsilon} \partial_x \left(\frac{1}{1+3x^2} \partial_x G(x, x_0) \right) dx = 1, \quad \forall \epsilon > 0.$$

This gives

$$\frac{1}{1+3x_0^2} (\partial_x G(x_0^+, x_0) - \partial_x G(x_0^-, x_0)) = 1,$$

i.e. $\partial_x G(x_0^+, x_0) = 1 + 3x_0^2 = c(1 + 3x_0^3)$. In conclusion $c = 1$ and $b = x_0 + x_0^3 - 2$. In other words,

$$G(x, x_0) = \begin{cases} x_0 + x_0^3 - 2 & \text{if } 0 \leq x < x_0 \\ x + x^3 - 2 & \text{if } x_0 < x \leq 1. \end{cases}$$

Question 2: Consider the operator $L : \phi \mapsto -\partial_x(x^{\frac{1}{2}}\partial_x\phi(x)) - \frac{\pi^2}{4}x^{-\frac{1}{2}}\phi(x)$, with domain $D = \{v \in \mathcal{C}^2(1,4); v(1) = 0, v(4) = 0\}$.

(i) What is the Null space of L ? (Hint: The general solution to $-\partial_x(x^{\frac{1}{2}}\partial_x\phi(x)) - \lambda x^{-\frac{1}{2}}\phi(x) = 0$ is $\phi(x) = c_1 \cos(2\sqrt{x}\sqrt{\lambda}) + c_2 \sin(2\sqrt{x}\sqrt{\lambda})$ for all $\lambda \geq 0$.)

Let ϕ be a member of the null space of L , say $\mathbf{N}(L)$. Then

$$-\partial_x(x^{\frac{1}{2}}\partial_x\phi(x)) - \frac{\pi^2}{4}x^{-\frac{1}{2}}\phi(x) = 0.$$

In other words, using the hint, $\phi(x) = c_1 \cos(2\sqrt{x}\sqrt{\lambda}) + c_2 \sin(2\sqrt{x}\sqrt{\lambda})$ with $\lambda = \frac{\pi^2}{4}$. The boundary conditions imply that

$$\phi(1) = 0 = -c_1, \quad \text{and} \quad \phi(4) = 0 = c_2 \sin(2\pi).$$

In conclusion $\mathbf{N}(L) = \text{span}\{\sin(\pi\sqrt{x})\}$, i.e., $\mathbf{N}(L)$ is a one-dimensional vector space.

(ii) Consider the problem $-\partial_x(x^{\frac{1}{2}}\partial_x\phi(x)) - \frac{\pi^2}{4}x^{-\frac{1}{2}}\phi(x) = \frac{1}{2}x^{-\frac{1}{2}}$, $x \in (1,4)$, with $\phi(1) = 0$, $\phi(4) = 0$. Does this problem have a solution? (Hint: $d(x^{\frac{1}{2}}) = \frac{1}{2}x^{-\frac{1}{2}}dx$.)

We are in the second case of the Fredholm alternative, since the null space of the operator L is not reduced to $\{0\}$. We must verify that $\frac{1}{2}x^{-\frac{1}{2}}$ is orthogonal to $\sin(\pi\sqrt{x})$. Using the hint and the change of variable $x^{\frac{1}{2}} = z$, we have

$$\int_1^4 \sin(\pi\sqrt{x}) \frac{1}{2}x^{-\frac{1}{2}} dx = \int_1^4 \sin(\pi x^{\frac{1}{2}}) d(x^{\frac{1}{2}}) = \int_1^2 \sin(\pi z) dz = -\frac{1}{\pi} [\cos(\pi z)]_1^2 = -\frac{2}{\pi}.$$

Hence $\int_1^4 \sin(\pi\sqrt{x}) \frac{1}{2}x^{-\frac{1}{2}} dx \neq 0$, which means that the above problem does not have a solution.

Question 3: Solve the following integral equation (Hint: $x^2 - 3xa + 2a^2 = (x - a)(x - 2a)$):

$$\int_{-\infty}^{+\infty} (f(y) - 3\sqrt{2}e^{-\frac{y^2}{2\pi}})f(x-y)dy = -4\pi e^{-\frac{x^2}{4\pi}}. \quad \forall x \in \mathbb{R}.$$

This equation can be re-written using the convolution operator:

$$f * f - 3\sqrt{2}e^{-\frac{x^2}{2\pi}} * f = -4\pi e^{-\frac{x^2}{4\pi}}.$$

We take the Fourier transform and use (4) to obtain

$$\begin{aligned} 2\pi\mathcal{F}(f)^2 - 2\pi 3\sqrt{2}\mathcal{F}(f) \frac{1}{\sqrt{4\pi\frac{1}{2\pi}}} e^{-\omega^2 \frac{1}{4\frac{1}{2\pi}}} &= -4\pi \frac{1}{\sqrt{4\pi\frac{1}{4\pi}}} e^{-\omega^2 \frac{1}{4\frac{1}{4\pi}}} \\ \mathcal{F}(f)^2 - 3\mathcal{F}(f)e^{-\omega^2 \frac{\pi}{2}} + 2e^{-\omega^2 \pi} &= 0 \\ (\mathcal{F}(f) - e^{-\omega^2 \frac{\pi}{2}})(\mathcal{F}(f) - 2e^{-\omega^2 \frac{\pi}{2}}) &= 0. \end{aligned}$$

This implies

$$\text{either } \mathcal{F}(f) = e^{-\omega^2 \frac{\pi}{2}}, \quad \text{or } \mathcal{F}(f) = 2e^{-\omega^2 \frac{\pi}{2}}.$$

Taking the inverse Fourier transform, we obtain

$$\text{either } f(x) = \sqrt{2}e^{-\frac{x^2}{2\pi}}, \quad \text{or } f(x) = 2\sqrt{2}e^{-\frac{x^2}{2\pi}}.$$

Question 4: Use the Fourier transform technique to solve the following PDE:

$$\partial_t u(x, t) + c \partial_x u(x, t) + \gamma u(x, t) = 0,$$

for all $x \in (-\infty, +\infty)$, $t > 0$, with $u(x, 0) = u_0(x)$ for all $x \in (-\infty, +\infty)$.

By taking the Fourier transform of the PDE, one obtains

$$\partial_t \mathcal{F}(u) - i\omega c \mathcal{F}(u) + \gamma \mathcal{F}(u) = 0.$$

The solution is

$$\mathcal{F}(u)(\omega, t) = c(\omega) e^{i\omega c t - \gamma t}.$$

The initial condition implies that $c(\omega) = \mathcal{F}(u_0)(\omega)$:

$$\mathcal{F}(u)(\omega, t) = \mathcal{F}(u_0)(\omega) e^{i\omega c t - \gamma t}.$$

The shift lemma in turn implies that

$$\mathcal{F}(u)(\omega, t) = \mathcal{F}(u_0(x - ct))(\omega) e^{-\gamma t} = \mathcal{F}(u_0(x - ct) e^{-\gamma t})(\omega).$$

Applying the inverse Fourier transform gives:

$$u(x, t) = u_0(x - ct) e^{-\gamma t}.$$

Question 5: Solve the following PDE by the method of characteristics:

$$\begin{aligned}\partial_t w + 3\partial_x w &= 0, & x > 0, & t > 0 \\ w(x, 0) &= f(x), & x > 0, & \text{ and } w(0, t) = h(t), & t > 0.\end{aligned}$$

First we parameterize the boundary of Ω by setting $\Gamma = \{x = x_\Gamma(s), t = t_\Gamma(s); s \in \mathbb{R}\}$ with

$$x_\Gamma(s) = \begin{cases} 0 & \text{if } s < 0, \\ s, & \text{if } s \geq 0. \end{cases} \quad \text{and} \quad t_\Gamma(s) = \begin{cases} -s & \text{if } s < 0, \\ 0, & \text{if } s \geq 0. \end{cases}$$

Then we define the characteristics by

$$\partial_t X(s, t) = 3, \quad \text{with} \quad X(s, t_\Gamma(s)) = x_\Gamma(s).$$

The general solution is $X(s, t) = 3(t - t_\Gamma(s)) + x_\Gamma(s)$. Now we make the change of variable $\phi(s, t) = w(X(s, t), t)$ and we compute $\partial_t \phi(s, t)$,

$$\partial_t \phi(s, t) = \partial_t w(X(s, t), t) + \partial_x w(X(s, t), t) \partial_t X(s, t) = \partial_t w(X(s, t), t) + 3\partial_x w(X(s, t), t) = 0.$$

This means that $\phi(s, t) = \phi(s, t_\Gamma(s))$. In other words

$$w(X(s, t), t) = w(X(s, t_\Gamma(s)), t_\Gamma(s)) = w(x_\Gamma(s), t_\Gamma(s)).$$

Case 1: If $s < 0$, then $X(s, t) = 3(t - t_\Gamma(s))$. This implies $t_\Gamma(s) = t - X/3$. The condition $s < 0$ and the definition $t_\Gamma(s) = -s$ imply $t - X/3 \geq 0$. Moreover we have

$$w(X, t) = w(0, t_\Gamma(s)) = h(t_\Gamma(s)).$$

In conclusion

$$w(X, t) = h(t - X/3), \quad \text{if} \quad 3t > X.$$

Case 2: If $s \geq 0$, then $X(s, t) = 3t + x_\Gamma(s)$. This implies $x_\Gamma(s) = X - 3t$. The condition $s \geq 0$ and the definition $x_\Gamma(s) = s$ imply $X - 3t \geq 0$. Moreover we have

$$w(X, t) = w(x_\Gamma(s), 0) = f(x_\Gamma(s)).$$

In conclusion

$$w(X, t) = f(X - 3t), \quad \text{if} \quad X \geq 3t.$$