

We are asked to solve  $\frac{\partial u}{\partial t} = k_1 \frac{\partial^2 u}{\partial x^2} + k_2 \frac{\partial^2 u}{\partial y^2}$  subject to the initial condition  $u(x, y, 0) = f(x, y)$ .

We begin by redefining the PDE as a spatial double Fourier transform:

$$\frac{\partial \bar{U}}{\partial t} = -(k_1 \omega_1^2 + k_2 \omega_2^2) \bar{U} \quad (1)$$

$$\bar{U}(\boldsymbol{\omega}, t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} u(x, y, t) e^{i\boldsymbol{\omega} \cdot \mathbf{r}} dx dy \quad (2)$$

where  $\boldsymbol{\omega}$  and  $\mathbf{r}$  are the wave number vector and position vector, respectively. Additionally we will define the “ $\mathbf{k}$ ” vector:

$$\boldsymbol{\omega} = \omega_1 \hat{i} + \omega_2 \hat{j} \quad (3)$$

$$\mathbf{r} = x \hat{i} + y \hat{j} \quad (4)$$

$$\mathbf{k} = k_1 \hat{i} + k_2 \hat{j} \quad (5)$$

The “elementary” solution, according to Haberman, can then be presented as:

$$\bar{U}(\boldsymbol{\omega}, t) = A(\boldsymbol{\omega}) e^{-\mathbf{k} \cdot \boldsymbol{\omega}^2 t} \quad (6)$$

$$A(\boldsymbol{\omega}) = \bar{U}(\boldsymbol{\omega}, 0) = \left(\frac{1}{2\pi}\right)^2 \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x, y) e^{i\boldsymbol{\omega} \cdot \mathbf{r}} dx dy \quad (7)$$

Thus, we can find a generic solution to this heat equation:

$$u(x, y, t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} A(\boldsymbol{\omega}) e^{-\mathbf{k} \cdot \boldsymbol{\omega}^2 t} e^{-i\boldsymbol{\omega} \cdot \mathbf{r}} d\omega_1 d\omega_2 \quad (8)$$

The convolution theorem is then applied, using the fact that  $H(\boldsymbol{\omega}) = F(\boldsymbol{\omega})G(\boldsymbol{\omega})$ .

$$h(\mathbf{r}) = \left(\frac{1}{2\pi}\right)^2 \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(\mathbf{r}_0) g(\mathbf{r} - \mathbf{r}_0) dx_0 dy_0 \quad (9)$$

With the knowledge that our solution is the product of  $A(\boldsymbol{\omega})$  and  $e^{-\mathbf{k} \cdot \boldsymbol{\omega}^2 t}$ , we can present a solution with the latter equation’s double Fourier transform.

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\mathbf{k} \cdot \boldsymbol{\omega}^2 t} e^{-i\boldsymbol{\omega} \cdot \mathbf{r}} d\omega_1 d\omega_2 \quad (10)$$

$$= \int_{-\infty}^{\infty} e^{-k_1 \omega_1^2 t} e^{-i\omega_1 x} d\omega_1 \int_{-\infty}^{\infty} e^{-k_2 \omega_2^2 t} e^{-i\omega_2 y} d\omega_2 \quad (11)$$

$$= \sqrt{\frac{\pi}{k_1 t}} e^{-x^2/4k_1 t} \sqrt{\frac{\pi}{k_2 t}} e^{-y^2/4k_2 t} \quad (12)$$

Equation (12) corresponds with our Green’s function,  $g(\mathbf{r})$ . Thus, because  $h(\mathbf{r})$  corresponds to our solution,  $u(\mathbf{r})$ , we can write:

$$h(x, y, t; x_0, y_0, 0) = u(x, y, t) = \left(\frac{1}{2\pi}\right)^2 \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(x_0, y_0) \frac{\pi}{t\sqrt{k_1 k_2}} \exp\left[-\frac{(x-x_0)^2}{4tk_1} - \frac{(y-y_0)^2}{4tk_2}\right] dx_0 dy_0 \quad (13)$$

Thus, our final solution is:

$$\mathbf{u(x, y, t)} = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \mathbf{f(x_0, y_0)} \frac{\mathbf{1}}{\mathbf{4t\pi\sqrt{k_1 k_2}}} \mathbf{exp}\left[-\frac{(x-x_0)^2}{\mathbf{4tk_1}} - \frac{(y-y_0)^2}{\mathbf{4tk_2}}\right] \mathbf{dx_0 dy_0} \quad (14)$$