

Test C – Solutions

Calculators may be used for simple arithmetic operations only!

Useful information:

Laplacian operator in polar coordinates:

$$\nabla^2 u = \frac{\partial^2 u}{\partial r^2} + \frac{1}{r} \frac{\partial u}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2}.$$

Bessel's equation:

$$\frac{\partial^2 Z}{\partial z^2} + \frac{1}{z} \frac{\partial Z}{\partial z} + \left(1 - \frac{n^2}{z^2}\right) Z = 0 \quad \text{has solutions } J_n(z) \text{ and } Y_n(z).$$

1. (60 pts.) Consider the eigenvalue problem

$$\phi'' = -\lambda\phi \quad (0 < x < 1); \quad \phi'(0) = \beta\phi(0), \quad \phi(1) = 0.$$

NOTE: You should be able to answer (c) and (d) even if you don't have a clue about (a) and (b).

(a) Assuming for now that the eigenvalues λ are all positive, find the eigenfunctions and show how to find the eigenvalues approximately (by a graphical method) if $\beta = 2$.

The easiest way is to write $\phi(x) = \sin[\omega(1-x)]$ (where $\omega^2 = \lambda$) to satisfy the second boundary condition, and then to calculate from the other boundary condition that $-\omega \cos \omega = \beta \sin \omega$. Another way is to write $\phi(x) = A \cos(\omega x) + B \sin(\omega x)$ with A and B unknown. The two boundary conditions give two homogeneous linear equations for A and B . Setting the determinant of the system equal to 0 yields again $-\omega \cos \omega = \beta \sin \omega$. Then you can solve for A and B up to a constant, getting something equivalent to $C \sin[\omega(1-x)]$ after a trig identity.

In either case, you then solve $-\frac{\omega}{\beta} = \tan \omega$ with $\beta = 2$ graphically as in Fig. 5.8.1 of Haberman.

(b) Show that the eigenvalues are indeed all positive if $\beta > 0$. Describe how an eigenvalue might be negative if $\beta < 0$.

Approach 1: Observe that

$$\begin{aligned} \lambda \int_0^1 |\phi(x)|^2 dx &\equiv \lambda \langle \phi, \phi \rangle = \langle \phi, (-\phi'') \rangle = - \int_0^1 \phi(x)^* \phi''(x) dx \\ &= -\phi(1)^* \phi'(1) + \phi(0)^* \phi'(0) + \int_0^1 |\phi'(x)|^2 dx = 0 + \beta |\phi(0)|^2 + \int_0^1 |\phi'(x)|^2 dx. \end{aligned}$$

Divide by $\langle \phi, \phi \rangle$ ("Rayleigh quotient") to see that λ is a sum of obviously positive quantities, provided that $\beta > 0$. If $\beta < 0$ we can't draw that conclusion, but to see clearly whether there *must* be a negative eigenvalue in that case, we need to resort to the second approach.

Approach 2: Suppose that $\lambda = -\rho^2$. Then the solution of the ODE is $\phi(x) = \sinh[\rho(1-x)]$, and a calculation analogous to the one in (a) leads to

$$\tanh \rho = -\frac{\rho}{\beta}.$$

This has a (nontrivial) solution if and only if $-1 < \beta < 0$, as shown by drawing a graph like Fig. 5.8.3. (For negative β there are still infinitely many positive eigenvalues, as shown by Figs. 5.8.2 and 5.8.4.)

(c) Write out the orthogonality and completeness relations obeyed by the eigenfunctions. Given the eigenfunctions $\phi_n(x) = \sin[\omega_n(1-x)]$ from (a) (or just the notation ϕ_n !) you can define orthonormal eigenfunctions

$$\psi_n(x) = \frac{\phi_n(x)}{\|\phi_n\|}, \quad \|\phi_n\|^2 \equiv \int_0^1 |\phi_n(x)|^2 dx.$$

(The integral can be evaluated, but it is *not* the usual $\pi/2$ or $L/2$ of ordinary Fourier series.) Then the orthogonality and completeness relations are

$$\int_0^1 \psi_n(x)^* \psi_m(x) dx = \delta_{mn}, \quad \sum_{n=1}^{\infty} \psi_n(x) \psi_n(y)^* = \delta(x-y).$$

(In this case the complex conjugates are superfluous, since we chose the eigenfunctions real. If there is a negative eigenvalue, you need to include a term for its eigenfunction in the completeness relation (but I should have excluded that case in the problem statement).)

You can also write the equivalent relations in terms of the ϕ_n , including the necessary factors of $\|\phi_n\|$.

(d) Use the (normalized) eigenfunctions and eigenvalues to solve the heat problem

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \quad (0 < x < 1, t > 0),$$

$$\frac{\partial u}{\partial x}(0, t) = \beta u(0, t), \quad u(1, t) = 0, \quad u(x, 0) = f(x).$$

I'll use the normalized functions ψ_n to keep the formulas simple. In the usual way you get (with $\lambda_n = \omega_n^2$, still assuming $\beta > 0$)

$$u(x, t) = \sum_{n=1}^{\infty} C_n \psi_n(x) e^{-\lambda_n t}.$$

Then

$$f(x) = \sum_{n=1}^{\infty} C_n \psi_n(x),$$

so

$$C_n = \langle \psi_n, f \rangle = \int_0^1 \psi_n(x)^* f(x) dx.$$

If you use unnormalized eigenfunctions, you get $\int_0^1 |\phi_n(x)|^2 dx$ in the denominator (without the square root, since you have redefined C_n).

2. (40 pts.) Do **ONE** of the following ((A) **OR** (B)). (Sorry, no extra credit for doing both. **Clearly indicate which question you want graded.**)

(A) Solve Laplace's equation in a sector,

$$\nabla^2 u = 0 \quad \left(0 < r < 3, 0 < \theta < \frac{\pi}{3} \right),$$

$$\frac{\partial u}{\partial \theta}(r, 0) = 0 = \frac{\partial u}{\partial \theta}\left(r, \frac{\pi}{3}\right), \quad u(3, \theta) = g(\theta).$$

Writing $u = R\Theta$ yields

$$\Theta'' = -\mu^2 \Theta, \quad R'' + \frac{1}{r} R' - \frac{\mu^2}{r^2} R = 0.$$

From the homogeneous boundary conditions we get $\Theta(\theta) = \cos(3m\theta)$, $\mu = 3m$, $m = 0, 1, \dots$ ($L = \pi/3 \Rightarrow m\pi/L = 3m$, $2/L = 6/\pi$.) The two independent solutions of the R equation are $r^{\pm 3m}$, and possibly $\ln r$ in the case $m = 0$, but for regularity at $r = 0$ we choose the nonnegative exponents (and no log). Thus

$$u(r, \theta) = \sum_{m=0}^{\infty} a_m r^{3m} \cos(3m\theta).$$

Then

$$g(\theta) = \sum_{m=0}^{\infty} a_m 3^{3m} \cos(3m\theta).$$

So if $m \neq 0$,

$$a_m = 3^{-3m} \frac{6}{\pi} \int_0^{\pi/3} \cos(3m\theta) g(\theta) d\theta,$$

$$a_0 = \frac{3}{\pi} \int_0^{\pi/3} g(\theta) d\theta.$$

(B) Solve the wave equation in a disk,

$$\nabla^2 u = \frac{\partial^2 u}{\partial t^2} \quad (0 \leq r < 3, 0 \leq \theta < 2\pi, 0 < t < \infty),$$

$$u(3, \theta, t) = 0, \quad u(r, \theta, 0) = f(r, \theta), \quad \frac{\partial u}{\partial t}(r, \theta, 0) = 0.$$

[This is a special case of the "drum problem" done in the class notes, pp. 106–107 and 112–115.]